Exercise 1. The Motion of Vortices

Consider a single two-dimensional vortex of strength ω placed in the origin. It produces a 2D velocity field given by

$$\mathbf{u}(\mathbf{r}) = \frac{\omega}{r} \mathbf{e}_{\theta},$$

which can be rewritten in components as

$$u_x = -\frac{\omega y}{r^2}, \ u_y = \frac{\omega x}{r^2}.$$

In this problem we will study the motion of many such vortices in their own velocity field. Suppose that we have n vortices of the same strength ω at distinct positions $\mathbf{r}_i = (x_i, y_i)$.

(a) Explain why the velocity of the i^{th} vortex is given by

$$\dot{x}_i = -\omega \sum_{j \neq i} \frac{y_i - y_j}{(x_i - x_j)^2 + (y_i - y_j)^2}, \ \dot{y}_i = \omega \sum_{j \neq i} \frac{x_i - x_j}{(x_i - x_j)^2 + (y_i - y_j)^2}$$
(1)

Quite surprisingly, this system can be described by using the Hamiltonian formalism. For each vortex we now define a generalized coordinate $q_i = x_i$ and identify $p_i = y_i$ with its 'canonical momentum'.

(b) Show that the equations (1) are equivalent to the Hamilton's equations

$$\dot{q}_k = \frac{\partial H}{\partial p_k}, \ \dot{p}_k = -\frac{\partial H}{\partial q_k} \text{ with } H = -\frac{\omega}{4} \sum_{\substack{i,j\\i\neq j}} \log\left[(q_i - q_j)^2 + (p_i - p_j)^2\right].$$
 (2)

Hint: After differentiating H it is convenient to write the sum as

$$\frac{\partial H}{\partial q_k} = \sum_{j \neq i} \sum_{i} A_{ij} \delta_{ik} + \sum_{i \neq j} \sum_{j} B_{ij} \delta_{jk}$$

for some objects A_{ij} , B_{ij} and use the properties of the Kronecker δ_{ij} .

(c) Consider an infinitesimal rotation of the system given by

$$\delta q_i = -\epsilon p_i, \ \delta p_i = \epsilon q_i, \ i = 1, 2, ..., n$$
 (3)

and show that it does not change the Hamiltonian to the order $\mathcal{O}(\epsilon)$. What is the generating function $\tilde{G}(q_1,...,q_n,p_1,...,p_n)$ of such transformation?

Hint: You should find the expansion $\log(1 + \epsilon^2) = \epsilon^2 + \mathcal{O}(\epsilon^4)$ useful. Recall that the generating function is defined by $\delta p_i = -\epsilon \partial \tilde{G}/\partial q_i$, $\delta q_i = \epsilon \partial \tilde{G}/\partial p_i$.

(d)* Define new coordinates $Q_i(\epsilon) = q_i + \delta q_i$, $P_i(\epsilon) = p_i + \delta p_i$ and the transformed Hamiltonian $H(\epsilon) = H(Q_i(\epsilon), P_i(\epsilon))$. Use the result from c) to deduce the value of $dH/d\epsilon|_{\epsilon=0}$. Hence show that \tilde{G} is a conserved quantity, i.e. $d\tilde{G}/dt = [\tilde{G}, H] = 0$.

Exercise 2. One more exercise about Canonical transformations...

Given the transformation for a system of one degree of freedom

$$Q = q \cos \alpha - p \sin \alpha$$

$$P = q \sin \alpha + p \cos \alpha$$
(4)

- (a) show that it satisfies the symplectic condition for any value of the parameter α .
- (b) Find a generating function F for the transformation. Hint: you can, for example, chose a generating function which depends on Q and q, so that p = p(Q, q) and P = P(Q, q), and use the transformation equations:

$$p = \frac{\partial F}{\partial q}$$
 and $P = -\frac{\partial F}{\partial Q}$. (5)

- (c) What is the physical significance of the transformation for $\alpha = 0$? For $\alpha = \pi/2$? Does your generating function work for both of these cases?
- (d) * If not, can you find a generating function valid for the case where the one you have just found doesn't work?

 Hint: try with a generating function which depends on other variables, e.g. q and P.

Exercise 3. Conserved tensor for harmonic oscillator

We have seen that the Kepler problem features, in addition to the total energy and angular momentum, a conserved quantity known as Laplace-Runge-Lenz vector. It turns out that the existence of conserved Laplace-Runge-Lenz-like vectors are not specific to the Kepler potential. Instead, it is rather a feature of radial force problems.

As such it is not surprising that also the three-dimensional isotropic harmonic oscillator allows for such a conserved quantity, which in this case is the (six component) tensor

$$\mathbf{A}_{ij} = \frac{1}{\lambda} (p_i p_j + \lambda^2 r_i r_j), \qquad (6)$$

where $i, j \in \{x, y, z\}$.

(a) Consider a particle of mass m in the harmonic potential

$$V(\vec{r}) = \frac{1}{2}u\vec{r}^2, \qquad \vec{r} \in \mathbb{R}^3.$$
 (7)

Write down the Lagrangian and find the Hamiltonian from the Legendre transformation.

(b) Using Poisson brackets $[\cdot, \cdot]$ and a suitable choice of λ , show that the tensor \mathbf{A}_{ij} defined in Eq. (6) is conserved, i.e.,

$$[\mathbf{A}_{ij}, H] = 0, \qquad \forall i, j. \tag{8}$$

Hint. Although the problem has radial symmetry, note that it might be simpler to carry out the calculation in cartesian coordinates.

(c) The tensor \mathbf{A}_{ij} has six components, corresponding to six constants of motion. The energy and the angular momentum are additional constants of motion, because the harmonic oscillator is a central force problem. Hence there are more constants of motion than degrees of freedom. For this reason some of them must be interlinked.

In this context, relate

$$\operatorname{tr} \mathbf{A}_{ij} = \sum_{i} \mathbf{A}_{ii} \tag{9}$$

to another conserved quantity.