

Quantized Hamiltonian

$$\mathcal{H} = \sum_{\vec{k}} \hbar \omega_{\vec{k}} \left(\hat{b}_{\vec{k}}^\dagger \hat{b}_{\vec{k}} + \frac{1}{2} \right) = \sum_{\vec{k}} \hbar \omega_{\vec{k}} \left(\hat{n}_{\vec{k}} + \frac{1}{2} \right)$$

creation/annihilation operators

$$\omega_{\vec{k}} = c_l |\vec{k}| \quad |\vec{k}| \leq k_D$$

Debye
wavevector

$$[\hat{b}_{\vec{k}}, \hat{b}_{\vec{k}'}^\dagger] = \delta_{\vec{k}, \vec{k}'}$$

like bosonic particles

$$[\hat{b}_{\vec{k}}, \hat{b}_{\vec{k}'}] = [\hat{b}_{\vec{k}}^\dagger, \hat{b}_{\vec{k}'}^\dagger] = 0$$

Phonons

displacement field operator

$$\hat{\vec{u}}(\vec{r}) = \frac{1}{\sqrt{\Omega}} \sum_{\vec{k}} \vec{e}_{\vec{k}} \sqrt{\frac{\hbar}{2\rho_m \omega_{\vec{k}}}} \left[\hat{b}_{\vec{k}} e^{i\vec{k} \cdot \vec{r}} + \hat{b}_{\vec{k}}^\dagger e^{-i\vec{k} \cdot \vec{r}} \right]$$

Correlation function

$$\begin{aligned}
g(\vec{r} - \vec{r}') &= \langle \hat{u}(\vec{r}) \cdot \hat{u}(\vec{r}') \rangle - \langle \hat{u}(\vec{r}) \rangle \cdot \langle \hat{u}(\vec{r}') \rangle = \langle \vec{u}(\vec{r}) \cdot \vec{u}(\vec{r}') \rangle \\
&= \frac{1}{\Omega} \sum_{\vec{k}, \vec{k}'} \frac{\hbar \vec{e}_{\vec{k}} \cdot \vec{e}_{\vec{k}'}}{2\rho_m \sqrt{\omega_{\vec{k}} \omega_{\vec{k}'}}} \langle [\hat{b}_{\vec{k}} e^{i\vec{k} \cdot \vec{r}} + \hat{b}_{\vec{k}}^\dagger e^{-i\vec{k} \cdot \vec{r}}] [\hat{b}_{\vec{k}'} e^{i\vec{k}' \cdot \vec{r}'} + \hat{b}_{\vec{k}'}^\dagger e^{-i\vec{k}' \cdot \vec{r}'}] \rangle \\
&= \frac{\hbar}{(2\pi)^2 \rho_m c_l} \int dk \frac{\sin(k|\vec{r} - \vec{r}'|)}{|\vec{r} - \vec{r}'|} (1 + 2\langle \hat{n}_{\vec{k}} \rangle)
\end{aligned}$$

$$g(R) \approx \frac{k_B \Theta_D}{(2\pi)^2 \lambda k_D} \times \begin{cases} \frac{1}{R^2} & T \ll T^*(R) & \text{quantum} \\ \frac{\pi T k_D}{\Theta_D R} & T \gg T^*(R) & \text{thermal} \end{cases}$$

$$T^*(R) = \frac{\Theta_D}{2k_D R}$$

melting transition in 3D:

mean square displacement

$$\langle \vec{u}(\vec{r})^2 \rangle = \frac{\hbar}{(2\pi)^2 \rho_m c_l} \int_0^{k_D} dk k \coth \left(\frac{\beta \hbar c_l k}{2} \right) = \begin{cases} \frac{k_D k_B T}{2\pi^2 \lambda} & T \gg \Theta_D, \\ \frac{k_D k_B \Theta_D}{8\pi^2 \lambda} & T \ll \Theta_D, \end{cases}$$

Lindemann criterion

lattice stable for

$$\langle \vec{u}(\vec{r})^2 \rangle < L_m^2 a^2$$

$$L_m \sim 0.1$$



melting temperature $k_B T_m = 2\pi \lambda a^3 L_m^2$

quantum melting $k_B \Theta_D < 8\pi \lambda a^3 L_m^2$

see He $\left\{ \begin{array}{ll} \text{solid} & \text{at high pressure} \\ \text{liquid} & \text{at low pressure} \end{array} \right.$

Anharmonic elastic energy - phonon-phonon interaction

$$E'_{\text{el}} = \frac{C'}{3} \int d^3r (\vec{\nabla} \cdot \vec{u})^3 \quad \text{third-order anharmonic coupling}$$

quantization \rightarrow $\mathcal{H}'_{\text{el}} = \frac{C}{3} \left(\frac{\hbar}{2\rho_m c_l \Omega} \right)^{3/2} \sum_{\vec{k}, \vec{k}'} (|\vec{k}| |\vec{k}'| |\vec{k} + \vec{k}'|)^{1/2}$

$$\times \left\{ \hat{b}_{\vec{k}} \hat{b}_{\vec{k}'} \hat{b}_{-\vec{k}-\vec{k}'} + 3 \hat{b}_{\vec{k}} \hat{b}_{\vec{k}'} \hat{b}_{\vec{k}+\vec{k}'}^\dagger + 3 \hat{b}_{\vec{k}+\vec{k}'} \hat{b}_{\vec{k}}^\dagger \hat{b}_{\vec{k}'}^\dagger + \hat{b}_{\vec{k}}^\dagger \hat{b}_{\vec{k}'}^\dagger \hat{b}_{-\vec{k}-\vec{k}'}^\dagger \right\} .$$

