

**Exercise 10.1 Magnetostriction in a Spin-Dimer-Model**

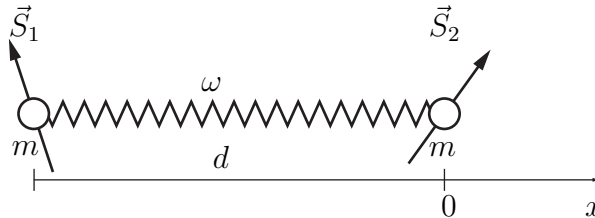
As in exercise 8.1, we again start with a dimer consisting of two (quantum) spins,  $s = 1/2$ , described by the Hamiltonian

$$\mathcal{H}_0 = J(\vec{S}_1 \cdot \vec{S}_2 + 3/4), \quad (1)$$

with  $J > 0$ . This time, however, the distance between the spins is not fixed but they are connected by a spring (cf. fig.) such that the Hamiltonian of the system reads

$$\mathcal{H} = \frac{\hat{p}^2}{2m} + \frac{m\omega^2}{2}\hat{x}^2 + J(1 - \lambda\hat{x})(\vec{S}_1 \cdot \vec{S}_2 + 3/4); \quad (2)$$

i.e., the spin-coupling constant depends on the distance between the two sites.



In the above figure,  $m$  is the mass of the two constituents,  $m\omega^2$  is the spring constant and  $d$  denotes the equilibrium distance between the two spins (in the case of no spin-spin interaction) from which the displacement  $x$  is measured.

- (a) Write the Hamiltonian (2) in second quantized form and calculate the partition sum, the internal energy, the specific heat and the entropy. In the limit  $T \rightarrow 0$ , discuss the entropy for different values of  $\lambda$ .

*Hint:* Set  $\hbar = 1$  and introduce an operator  $\hat{n}_t$  defined through

$$\langle \sigma | \hat{n}_t | \sigma \rangle = \begin{cases} 1 & \sigma \text{ is a triplet,} \\ 0 & \sigma \text{ is the singlet,} \end{cases} \quad (3)$$

where  $|\sigma\rangle$  denotes the spin-dependent part of the dimer state. Trace first over the spin-degrees of freedom.

- (b) Calculate the expectation value of the distance of the two spins,  $\langle d + \hat{x} \rangle$ , as well as the fluctuations  $\langle (d + \hat{x})^2 \rangle$ .

How are these quantities affected by a magnetic field in  $z$ -direction, i.e., by an additional term in (2) of the form

$$\mathcal{H}_m = -g\mu_B H \sum_{i,m} S_{i,m}^z? \quad (4)$$

- (c) If the two sites are oppositely charged, i.e.,  $\pm q$ , the dimer forms a dipole with moment  $P = q(d + x)$ . This dipole moment can be measured by applying an electric field  $E$  in  $x$ -direction,

$$\mathcal{H}_{el} = -q(d + \hat{x}) \cdot E. \quad (5)$$

Calculate the zero-field susceptibility of the dimer,

$$\chi_0^{(el)} = - \left. \frac{\partial^2 F}{\partial E^2} \right|_{E=0}, \quad (6)$$

and compare your result with the fluctuation-dissipation theorem which states

$$\chi_0^{(el)} \propto \left( \langle (d + \hat{x})^2 \rangle - \langle d + \hat{x} \rangle^2 \right). \quad (7)$$

Plot the zero-field susceptibility as a function of the applied magnetic field  $H$  and discuss your result.

### Exercise 10.2 The Ising Model in the High-Temperature Limit

Consider the Ising model with nearest neighbor interactions in the presence of a homogeneous magnetic field  $h$ ,

$$\mathcal{H} = -\frac{J}{2} \sum_{\langle i,j \rangle} S_i \cdot S_j - h \sum_i S_i, \quad J > 0, \quad (8)$$

where the spins assume the values  $S_i = \pm S$  along the magnetic field and the sum  $\sum_{\langle i,j \rangle}$  runs over nearest neighboring sites on the lattice. The number of spins is very large,  $N \gg 1$ , such that surface effects may be neglected.

- a) Determine the partition function in the high-temperature limit  $\beta J \ll 1$ .  
*Hint:* Note that for  $\beta J \ll 1$ , one may neglect bond-bond correlations and the partition function simplifies to

$$Z \approx \sum_{\{S_i\}} e^{\beta h \sum_i S_i} \left( 1 + \frac{\beta J}{2} \sum_j \sum_{m \in \Lambda_j} S_j \cdot S_m \right), \quad (9)$$

where  $\Lambda_j$  represents the set of nearest neighbors of site  $j$  such that  $|\Lambda_j| = z$  with the coordination number  $z$ .

- b) Calculate the spin susceptibility at  $h = 0$ . In analogy to the lecture notes (Section 3.4.6), plot  $1/\chi_0$  as a function of temperature and extrapolate the high- $T$  limit to lower temperatures to find the intersection on the  $T$ -axis. What is the physical interpretation of the intersection temperature?

**Office Hours:** Monday, November 23, 8-10 am (HIT K 43.2)