

Deformation of an elastic medium

when point force is applied in the bulk

Equilibrium equation is

$$\nabla^2 \vec{u} + \frac{1}{1-2\beta} \text{grad div } \vec{u} = - \frac{2(1+\beta)}{E} \vec{F} \delta(r)$$

point force

We should find Green's function.

It is convenient to do it in Fourier space.

$$\kappa^2 \vec{u} + \frac{1}{(1-2\beta)} \vec{\kappa} (\vec{\kappa} \cdot \vec{u}) = \frac{2(1+\beta)}{E} \vec{F}$$

To solve it we extract $\vec{\kappa} \cdot \vec{u}$

$$\kappa^2 (\vec{\kappa} \cdot \vec{u}) + \frac{\kappa^2 (\vec{\kappa} \cdot \vec{u})}{1-2\beta} = \frac{2(1+\beta)}{E} (\vec{\kappa} \cdot \vec{F})$$

Then


$$(\vec{k} \cdot \vec{u}) = \frac{(1+\beta)(1-2\beta)}{E(1-\beta)} \frac{\vec{k} \cdot \vec{F}}{k^2}$$

Substituting it back in the equation for u we obtain

$$\bullet u(\vec{k}) = \frac{2(1+\beta)}{E} \left[\frac{\vec{F}}{k^2} - \frac{1}{2(1-\beta)} \frac{\vec{k} (\vec{k} \cdot \vec{F})}{k^4} \right]$$

Calculating Fourier transform back

$$\int \frac{e^{-i\vec{k}\vec{r}}}{k^2} \frac{d^3k}{(2\pi)^3} = \frac{1}{(2\pi)^3} \int \frac{e^{-ik_z r}}{k_z^2 + k_\perp^2} dk_z d^2k_\perp =$$

$$\bullet = \frac{2\pi}{(2\pi)^3} \int \frac{e^{-k_\perp r}}{2k_\perp} d^2k_\perp$$


$$= \frac{2\pi^2}{(2\pi)^3} \int_0^\infty e^{-k_\perp r} \frac{\pi k_\perp dk_\perp}{k_\perp} =$$

$$= \frac{2\pi^2}{(2\pi)^3} \frac{1}{r} = \frac{1}{4\pi r}$$

(Also known from electrostatic)

Because Fourier transform of ∇f is $i\vec{k} f(k)$, Fourier transform of $\vec{k} (\vec{k} \cdot \vec{F}) f(k)$ is $-\nabla (\vec{F} \cdot \nabla) f(r)$ (56)

Then

$$\int e^{-i\vec{k} \cdot \vec{r}} \frac{\vec{k} \cdot (\vec{k} \cdot \vec{F})}{k^4} \frac{d^3 k}{(2\pi)^3} = -\vec{\nabla} (\vec{F} \cdot \nabla) \int \frac{e^{-i(\vec{k} \cdot \vec{r})}}{k^4} \frac{d^3 k}{(2\pi)^3}$$

The last integral is

$$\frac{1}{(2\pi)^3} \int \frac{e^{-ik_z r} dk_z d^2 k_\perp}{(k_z^2 + k_\perp^2)^2} = \frac{1}{(2\pi)^3} \int \frac{e^{-ik_z r} dk_z \pi d(k_\perp^2)}{(k_z^2 + k_\perp^2)^2} =$$

$$= \frac{\pi}{(2\pi)^3} \int \frac{e^{-ik_z r} dk_z}{k_z^2}$$

$$\frac{\partial^2 I}{\partial r^2} = - \int e^{-ik_z r} dk_z = -2\pi \delta(r) \Rightarrow$$

$$I = -\pi |r|$$

Combining all factors we obtain

$$\int e^{-i\vec{k}\cdot\vec{r}} \frac{\vec{k}\cdot(\vec{k}\cdot\vec{F})}{k^4} \frac{d^3k}{(2\pi)^3} = \frac{\pi^2}{(2\pi)^3} \vec{\nabla}\cdot(\vec{F}\cdot\vec{\nabla}) \frac{1}{r}$$

$$= \frac{1}{8\pi} \frac{\vec{F} - \vec{n}(\vec{n}\cdot\vec{F})}{r}, \text{ with } \vec{n} = \frac{\vec{r}}{r}$$

• Substituting back we arrive at

$$u = \frac{(1+\beta)}{2\pi E} \left[\frac{\vec{F}}{r} - \frac{1}{4(1-\beta)} \frac{\vec{F} - \vec{n}(\vec{n}\cdot\vec{F})}{r} \right]$$

$$\vec{u} = \frac{1+\beta}{8\pi E(1-\beta)} \frac{(3-4\beta)\vec{F} + \vec{n}(\vec{n}\cdot\vec{F})}{r}$$

• Note that displacement decays

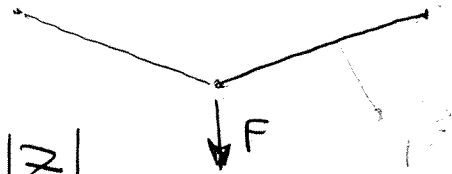
only as $\frac{1}{r}$. This is the property of Laplace equation in 3d

In 1 dimension (string)

$$F = \int \left[\frac{C}{2} \left(\frac{du}{dz} \right)^2 + F \delta(z) u \right] dz$$

Then equilibrium equation is

$$C \frac{d^2 u}{dz^2} = F \delta(z)$$



• It has solution $u = -d|z|$

Integrating around $z=0$ we obtain

$$C \left. \frac{du}{dz} \right|_{-0}^{+0} = F \quad \Rightarrow \quad d = \frac{F}{2C}$$

• In this case displacement from the action of the point force grows linearly with the distance.