

Elasticity of classical harmonic lattice

(25)

Consider lattice of atoms, that interact with pair potential $V(R - R')$

The total energy is

$$E_{\text{int}} = \frac{1}{2} \sum_{RR'} V(\vec{R} - \vec{R}')$$



- If atoms are displaced then the energy is

$$E_{\text{int}} = \frac{1}{2} \sum_{RR'} V(\vec{R} - \vec{R}' + \vec{u}(R) - \vec{u}(R')) =$$

$$= E_0 + \frac{1}{2} \sum_{R, R'} \frac{\partial V}{\partial R_i} [u_i(R) - u_i(R')] +$$

$$+ \frac{1}{4} \sum_{R, R'} \frac{\partial^2 V(R - R')}{\partial R_i \partial R_j} [u_i(R) - u_i(R')][u_j(R) - u_j(R')]$$

Because initial positions ($u = 0$) were equilibrium positions the linear in $u(R)$ term vanishes. As a result we obtain

$$\delta E_{\text{int}} = \frac{1}{4} \sum_{R, R', i, j} C_{ij}(R-R') [u_i(R) - u_i(R')] [u_j(R) - u_j(R')] \quad (26)$$

with $C_{ij} = \frac{\partial^2 V(R)}{\partial R_i \partial R_j}$

Expanding $u_i(R) = u_i(R') + \frac{\partial u_i}{\partial R_j} (R - R')$,

and shifting $R - R' \rightarrow R$ we rewrite

$$\bullet \quad \delta E_{\text{int}} = \frac{N}{4} \sum_R C_{ik}(R) R_j R_\ell \frac{\partial u_i}{\partial R_j} \frac{\partial u_k}{\partial R_\ell}$$

N is the total number of atoms

Since energy does not change under rotation we can replace

$$\bullet \quad \frac{\partial u_i}{\partial R_j} = \frac{1}{2} \left(\frac{\partial u_i}{\partial R_j} + \frac{\partial u_j}{\partial R_i} \right) + \frac{1}{2} \left(\cancel{\frac{\partial u_i}{\partial R_j}} - \cancel{\frac{\partial u_j}{\partial R_i}} \right) \rightarrow u_{ij}$$

Then

$$\delta E_{\text{int}} = \frac{N}{4} \sum_R C_{ik} R_j R_\ell u_{ij} u_{k\ell}$$

To write it in symmetrized form let us interchange indices and $i \leftrightarrow j, k \leftrightarrow l$ and add

$$\delta E_{int} = \frac{N}{16} \sum_{\mathbf{R}} [C_{ik} R_j R_e u_{ij} u_{ke} + C_{jk} R_i R_e u_{ji} u_{ki}] \quad (27)$$

$$+ C_{ie} R_j R_k u_{ij} u_{ek} + C_{je} R_i R_k u_{ji} u_{ek}$$

Since $u_{ij} = u_{ji}$, $u_{ke} = u_{ek}$
 all $u_{ij}, u_{ji}, u_{ke}, u_{ek}$ are the same and we
 obtain

$$\delta E_{int} = \frac{N}{16} \sum_{\mathbf{R}} [C_{ik} R_j R_e + C_{jk} R_i R_e + C_{ie} R_j R_k + C_{je} R_i R_k] u_{ij} u_{ke}$$

Elastic tensor is defined as

$$\delta E_{el} = \frac{1}{2} \int dV \lambda_{ijke} u_{ij} u_{ke} \Rightarrow$$

$$\lambda_{ijke} = \frac{1}{8V_0} \sum_{\mathbf{R}} [C_{ik}(R) R_j R_e + C_{jk}(R) R_i R_e +$$

$$+ C_{ie}(R) R_j R_k + C_{je}(R) R_i R_k]$$

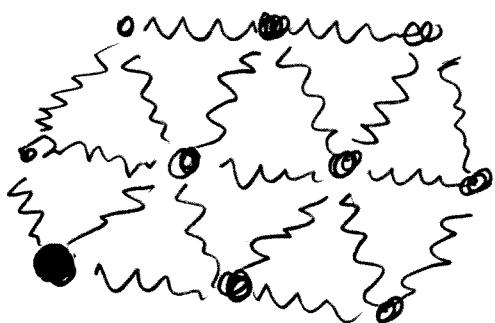
where V_0 is the unit cell volume

$$V = N V_0$$

Example

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Consider hexagonal lattice of mass points connected by harmonic springs between nearest neighbors.



Interaction potential

$$V = \frac{d}{2} (\sqrt{x^2 + y^2} - a)^2 =$$

$$= \frac{d}{2} (x^2 + y^2 + a^2 - 2a\sqrt{x^2 + y^2})$$

$$C_{xx} = V''_{xx} = d \left(1 - \frac{ay^2}{(x^2 + y^2)^{3/2}} \right) = d \left(1 - \frac{y^2}{a^2} \right)$$

$$C_{yy} = V''_{yy} = d \left(1 - \frac{x^2}{a^2} \right)$$

$$C_{xy} = V''_{xy} = d \frac{xy}{a^2}$$

$$\begin{aligned} \lambda_{ijk\epsilon} = & \frac{1}{8S_0} \sum_R [C_{ik}(R) R_j R_\epsilon + C_{jk}(R) R_i R_\epsilon + \\ & + C_{ie}(R) R_j R_k + C_{je}(R) R_i R_k] \end{aligned}$$

$$S_0 = \boxed{\text{square}} = a^2 \frac{\sqrt{3}}{2}$$

We should sum over 6 neighbors (29)

$$\lambda_{xxxx} = \frac{1}{8S_0} 4 \sum_R C_{xx}(R) x^2 = \frac{\alpha a^2}{8S_0} 4 \cdot \left(2 + 4 \cdot \frac{1}{4} \cdot \frac{1}{4} \right) = \frac{9 \alpha a^2}{8S_0}$$

$$\lambda_{xxyy} = \frac{1}{8S_0} 4 \sum_R C_{xy}(R) xy = \frac{\alpha a^2}{8S_0} 4 \sum \frac{x^2 y^2}{a^4} = \frac{3 \alpha a^2}{8S_0}$$

$$\begin{aligned} \lambda_{xyxy} &= \frac{1}{8S_0} \sum_R [C_{xx}(R) y^2 + 2C_{xy} xy + C_{yy} x^2] = \\ &= \frac{\alpha}{8S_0} \left[\frac{3}{4} + 2 \cdot \frac{3}{4} + \frac{3}{4} \right] = \frac{3 \alpha a^2}{8S_0} \end{aligned}$$

$$\begin{aligned} \text{Thus } E_{el} &= \frac{\lambda_{xxxx}}{2} (u_{xx}^2 + u_{yy}^2) + \lambda_{xxyy} u_{xx} u_{yy} + \\ &\quad + 2 \lambda_{xyxy} u_{xy}^2 = \end{aligned}$$

$$= \frac{3 \alpha a^2}{16S_0} [3(u_{xx}^2 + u_{yy}^2) + 2 u_{xx} u_{yy} + 4 u_{xy}^2] =$$

$$\begin{aligned} &= \frac{3 \alpha a^2}{8S_0} \left[\underbrace{(u_{xx} + u_{yy})^2}_{\frac{1}{2} U_{ee}^2} + \underbrace{(u_{xx}^2 + u_{yy}^2 + 2 u_{xy}^2)}_{M U_{ik}^2} \right] \Rightarrow M = 2 \lambda = \frac{3 \alpha a^2}{8S_0} \\ &\quad K = \lambda + M = \frac{3}{2} M \end{aligned}$$