

## Corbino disk and the integer quantum Hall effect

## Exercise 9.1 The lowest Landau level in the Corbino geometry

The Hamilton operator for an electron ( $e < 0$ ) restricted to the plane  $z = 0$  and exposed to a magnetic field is given by

$$H = \frac{1}{2m^*} \left( \mathbf{p} - \frac{e}{c} \mathbf{A} \right)^2 + U(r) \quad (1)$$

where  $r^2 = x^2 + y^2$ . The annular potential

$$U(r) = \frac{C_1}{r^2} + C_2 r^2 + C_3 \quad (2)$$

with  $C_1, C_2 > 0$  yields a Corbino<sup>1</sup> geometry confining the electron on a two-dimensional ring. The constant term  $C_3$  only leads to a shift of the energy which is why we neglect it in the following. Let the magnetic field  $\mathbf{B}$  be homogeneous and directed along the  $z$ -axis for  $r > 0$ . In addition, a magnetic flux  $\Phi = \nu \Phi_0$  through the origin ( $r = 0$ ) that does not physically touch the electron is assumed:

$$\mathbf{B} = [B + \nu \Phi_0 \delta(\mathbf{r})] \mathbf{e}_z, \quad (3)$$

with  $B, \nu > 0$ .  $\Phi_0 = hc/|e| = 2\pi\hbar c/|e|$  is the magnetic flux quantum.

a) Show that the vector potential can be chosen in the symmetric gauge

$$\mathbf{A} = \frac{1}{2} \left( B + \frac{\nu \Phi_0}{\pi r^2} \right) (x \mathbf{e}_y - y \mathbf{e}_x). \quad (4)$$

b) Let us now solve this single-particle problem in the symmetric gauge. To begin with we consider  $\nu = 0$  and use the following ansatz for the wave functions of the lowest Landau level:

$$\psi_m(r, \phi) = A r^\alpha e^{-im\phi} e^{-\frac{r^2}{4l^{*2}}}. \quad (5)$$

In order to avoid a singularity at the origin we have to demand that  $\alpha \geq 0$ . Furthermore, the uniqueness of the wave function is ensured by  $m \in \mathbb{Z}$ . Show that the following Schrödinger equation for the radial part is obtained:

$$\left\{ \frac{\hbar^2}{2m^*} \left[ -\frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial}{\partial r} + \left( \frac{m}{r} - \frac{r}{2l^2} \right)^2 \right] + U(r) - E_m \right\} r^\alpha e^{-\frac{r^2}{4l^{*2}}} = 0, \quad (6)$$

where  $l^2 = \hbar c/|eB|$ .

**Hint:**  $L_z = xp_y - yp_x = -i\hbar \partial_\phi$  and  $\partial_x^2 + \partial_y^2 = (1/r) \partial_r r \partial_r + (1/r)^2 \partial_\phi^2$ .

Show that the following relations are obtained

$$\alpha = \sqrt{m^2 + C_1^*}, \quad \frac{1}{l^{*2}} = \frac{1}{l^2} \sqrt{1 + C_2^*}, \quad (7)$$

<sup>1</sup>After the Italian physicist O. M. Corbino (1876-1937).

where  $C_1^* = 2m^*C_1/\hbar^2$  and  $C_2^* = 8l^4m^*C_2/\hbar^2$  are dimensionless parameters. Furthermore, show that the energy is given by

$$E_m = \frac{\hbar\omega_c}{2} \left[ \frac{l^2}{l^{*2}}(\alpha + 1) - m \right] \quad (8)$$

where  $\omega_c = |eB|/m^*c$ .

- c) (i) Consider the case  $U \equiv 0$ . Plot the radial part of the wave function for several  $m$ 's and show that it has a maximum at  $r_m = \sqrt{2ml}$ . Thus, the wave functions are localized on circles of radius  $r_m$ . Compute the magnetic flux penetrating the circle of radius  $r_m$ . How big is the degeneracy of the lowest Landau level in this case?

(ii) Show that for  $C_1^*, C_2^* \ll 1$  and  $m \gg 1$  the energy Eq. (8) can be approximated by

$$E_m \approx \hbar\omega_c/2 + U(r_m).$$

Thus, the wave functions are localized on curves of equal potential energy.

- d) Compute the angular speed  $\omega_m$  in the state  $\psi_m$

$$\omega_m := \langle \psi_m | v_\phi / r | \psi_m \rangle = -\frac{1}{\hbar} \frac{\partial E_m}{\partial m} \quad (9)$$

and show that under the conditions stated in c) (ii) one finds

$$\omega_m \approx -\frac{U'(r_m)}{m^*\omega_c r_m}. \quad (10)$$

Why are the states which are located on different boundaries of the Corbino-disk called *chiral* edge states?

- e) We assume now  $\nu > 0$  and use the same ansatz as above for the wave functions of the lowest Landau level. Show that the resulting Schrödinger equation is equal to Eq. (6) except for the substitution  $m \rightarrow m - \nu$ .

Use gauge transformations of the vector potential to show that under certain conditions the contribution from the magnetic flux through the origin can be transformed away. What is the condition on  $\nu$ ?

- f) What happens if the magnetic flux through the central hole is turned on adiabatically from  $\nu = 0$  to  $\nu = 1$ ? Use the above findings to recapitulate Laughlin's gauge argument for the quantization of the Hall conductivity (see chapter 4.2.2). What is the role of the additional magnetic flux?

**Office hour:**

Monday, May 2, 2011 - 9:00 to 11:00 am

HIT K 31.3

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