

**Exercise 13.1 Critical temperature in the Stoner model**

We consider three types of dispersion relations:

- $\epsilon_{\vec{k}} = \epsilon_0 \pm \frac{\hbar^2 \vec{k}^2}{2m}$  (3D) and
- $\epsilon_k = \epsilon_0 + \alpha k$  (1D).

Plot the critical temperature of the Stoner model for fixed interaction strength  $U$  depending on the chemical potential  $\mu$ .

**Exercise 13.2 Stoner instability**

In the lecture, it was shown that a system described by the mean-field Hamiltonian

$$\mathcal{H}_{\text{MF}} = \frac{1}{\Omega} \sum_{\vec{k}, s} (\epsilon_{\vec{k}} + U n_{-s}) c_{\vec{k}s}^\dagger c_{\vec{k}s} - U n_\uparrow n_\downarrow \quad (1)$$

shows an instability towards a magnetically ordered state at  $N(\epsilon_F)U_C = 1$  (note that here,  $N(\epsilon)$  is the density of states per spin).

Show for the case of a parabolic dispersion and  $T = 0$  that there are actually three distinct states:

- a paramagnetic state:  $N(\epsilon_F)U < 1$ ,
- an imperfect ferromagnetic state:  $3/2^{4/3} > N(\epsilon_F)U > 1$  and
- a perfect ferromagnetic state:  $N(\epsilon_F)U > 3/2^{4/3}$ .

**Hint:** Introduce a variable for the magnitude of the polarization

$$\frac{N_\uparrow}{N_e} = \frac{1}{2}(1+x) \quad \frac{N_\downarrow}{N_e} = \frac{1}{2}(1-x) \quad (2)$$

where  $N_{\uparrow(\downarrow)}$  is the total number of up-spins (down-spins) and  $N_e$  is the total number of electrons. Minimize the total energy of the system with respect to  $x$ .

Plot the polarization of the system  $x$  as a function of  $N(\epsilon_F)U$ .

**Exercise 13.3 Particle-Hole Excitations in Itinerant Ferromagnets**

In section 6.3 of the lecture notes the low-energy spectrum of (magnons) spin-waves in itinerant ferromagnets was derived. It is crucial for the existence of well-defined magnons that elementary particle-hole excitations are gapped. Try to explain (without detailed calculations) why there is such a gap and why it is important for the observability of magnons!